

# Dislocation mechanics of high rate deformations

Ron Armstrong\* (and Qizhen Li\*\*)

\*University of Maryland, College Park, MD 20742

\*\* University of Nevada at Reno, NV 89557

1. Dislocation dynamics:  $(d\varepsilon/dt) = (1/m)\rho bv \rightarrow (1/m)(dp/dt)b\Delta x_d$  8 charts
  - 1.a. TASRA, Zerilli-Armstrong (Z-A) and Johnson-Cook relations
  - 1.b. The dislocation activation volume,  $v^* = W_0/\tau_{Th}$
  - 1.c. Application to copper, steel and tantalum
2. Shear banding (1981/2, 1993/4) 4 charts
  - 2.a. dislocation pile-up avalanche:  $n\tau = \tau^*$
  - 2.b. Lueders-type and other type shear banding (Fe, Ti, Mg)
3. Shock loading *versus* isentropic compression experiments 6 charts
  - 3.a. Shock-induced dislocation generations:  $(d\varepsilon/dt) = (1/m)(dp/dt)b\Delta x_d$
  - 3.b. Control by dislocation drag in ICEs:  $\sigma_{Th} = (c/\beta_1)(d\varepsilon/dt)$
4. Hall-Petch for nanopolycrystals:  $\sigma_\varepsilon = \sigma_{0\varepsilon} + k_\varepsilon l^{-1/2}$  5 charts
  - 4.a. The dislocation number dependence:  $\Delta n = 1.0$ ;  $\Delta\tau = (-1/n)\tau$
  - 4.b. H-P strain rate sensitivity:  $v^{*-1} = v_0^{*-1} + (k_\varepsilon/2m\tau_{Ce}v_C^*)l^{-1/2}$

# Dislocation Dynamics: the TASRA

(1) Johnson - Cook:  $\sigma = (A + B\varepsilon^n)(1 + C\ln[d\varepsilon/dt])(1 - T^m)$   
 $T^* = (T - T_A)/(T - T_M)$

(2) Z - A:  $[d\gamma/dt] = [d\gamma/dt]_0 \exp\{(-G_0 + \int v^* d\tau_{Th})/kT\}$   
 $v^* = kT \{\partial \ln[d\gamma/dt]/\partial \tau_{Th}\}_T = W_0/T_{Th}$

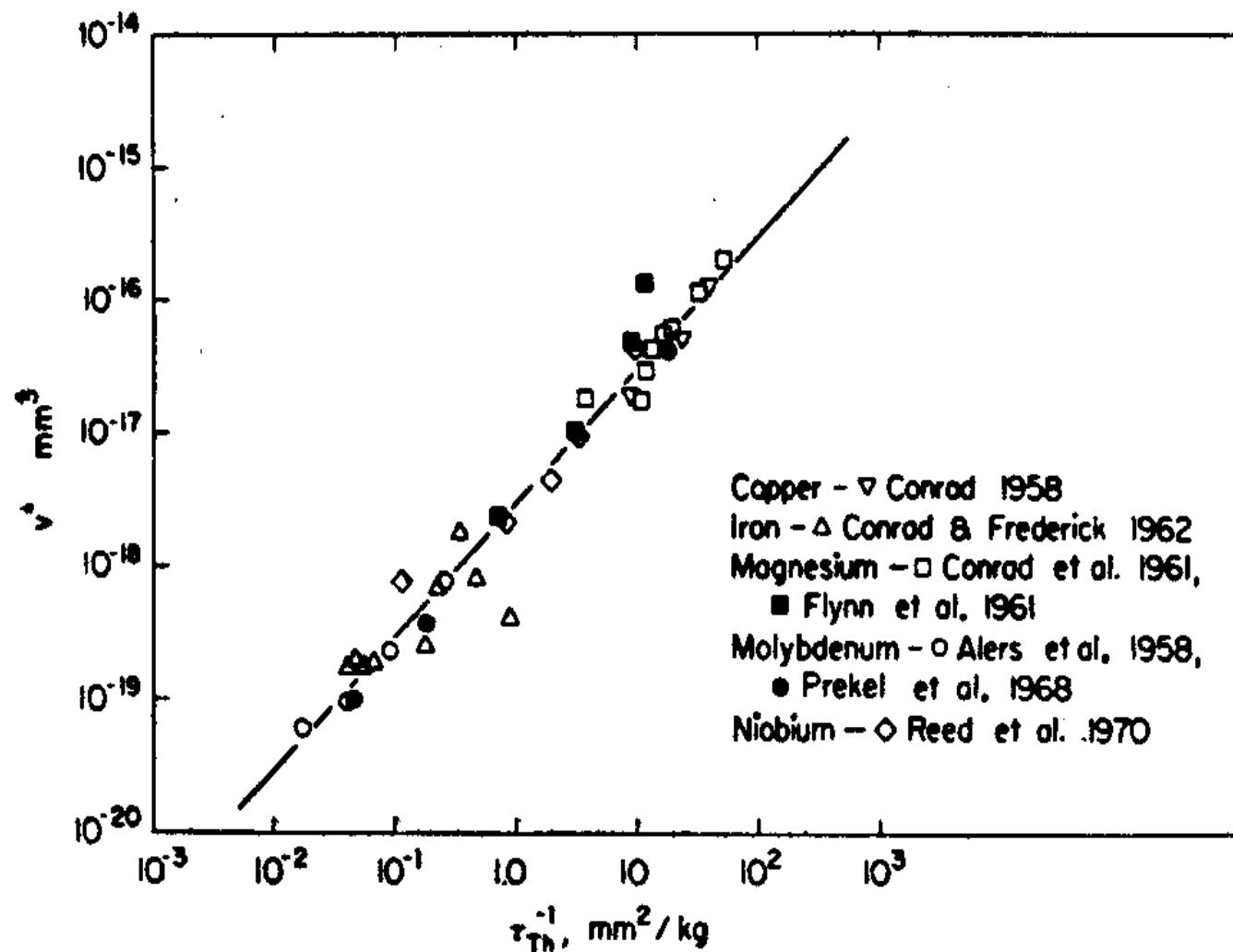
(2.1) bcc case:  $\sigma_\varepsilon = m\tau_\varepsilon = \sigma_{0G} + B_0 \exp\{-\beta T\} + K\varepsilon^n + k_\varepsilon l^{-1/2}$

(2.2) fcc:  $\sigma_\varepsilon = m\tau_\varepsilon = \sigma_{0G} + B_1 [\varepsilon_r \{1 - \varepsilon/\varepsilon_r\}]^{1/2} \exp\{-\beta T\} + k_\varepsilon l^{-1/2}$   
 $\beta = \beta_0 - \beta_1 \ln[d\varepsilon/dt]$

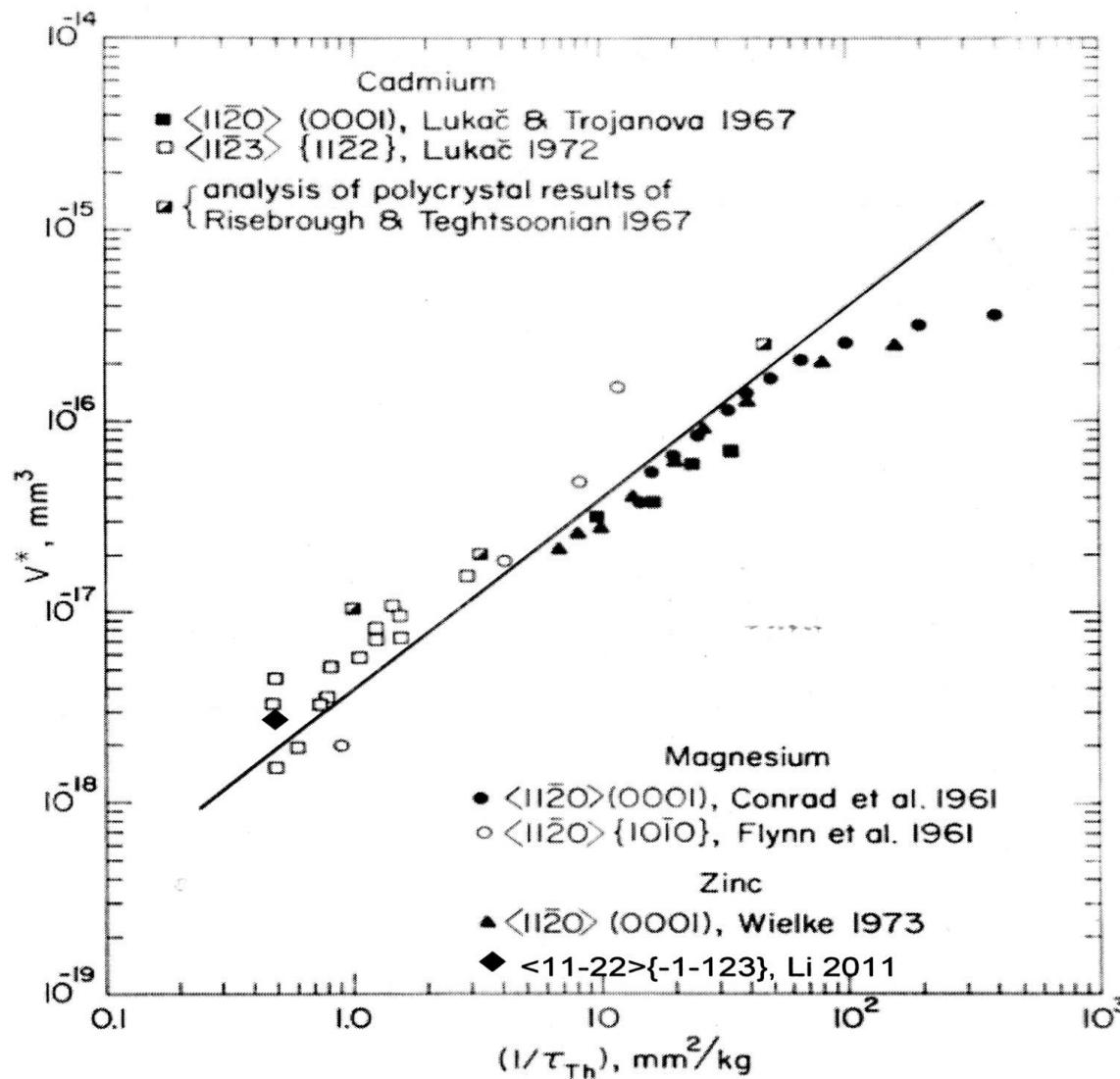
(3) (Hall-Petch) Twinning:  $\sigma_T = \sigma_{0T} + k_T l^{-1/2}; k_T > k_\varepsilon$

R.W. Armstrong, "Thermal Activation – Strain Rate Analysis (TASRA) for Polycrystalline Metals", (*Indian J. Scient. Indust. Res.*, **32**, 591-598 (1973); F.J. Zerilli and R.W. Armstrong, Dislocation-mechanics-based constitutive relations for material dynamics calculations, *J. Appl. Phys.* **61**, 1816-1824 (1987); R.W. Armstrong, "Dislocation mechanics description of polycrystal plastic flow and fracturing", In: *Mechanics and Materials: Fundamentals and Linkages*, M.A. Meyers, R.W. Armstrong and H.O.K. Kirchner (John Wiley & Sons, Inc., NY, 1999), pp. 363-398

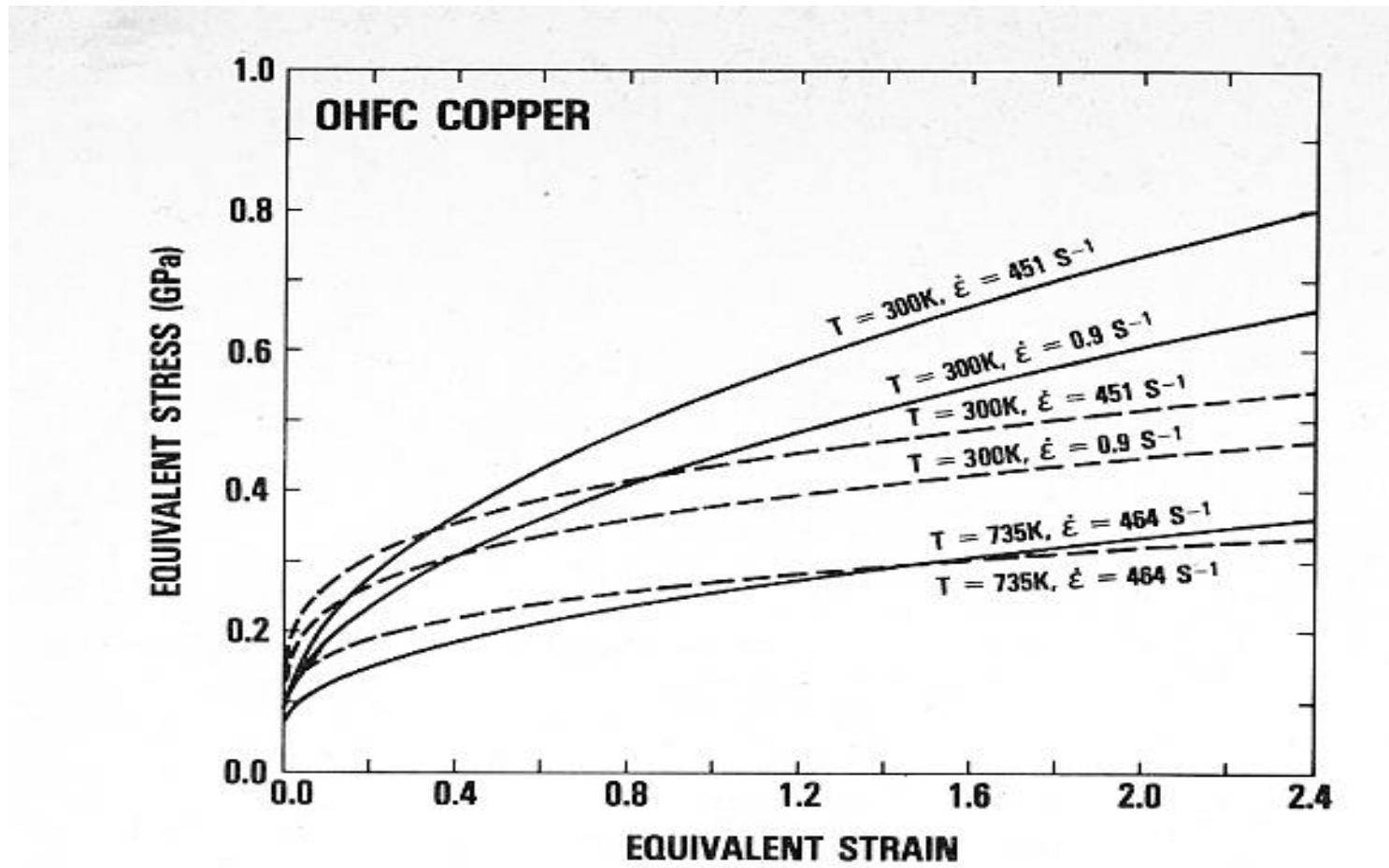
# A critical constitutive equation role for the thermal activation volume, $v^* = bA^* = kT \{\partial \ln[d\gamma/dt]/\partial T_{Th}\}_T$



# Activation volume, $v^*$ , measurements for hcp metals



# Z-A and J-C stress-strain curves for Cu

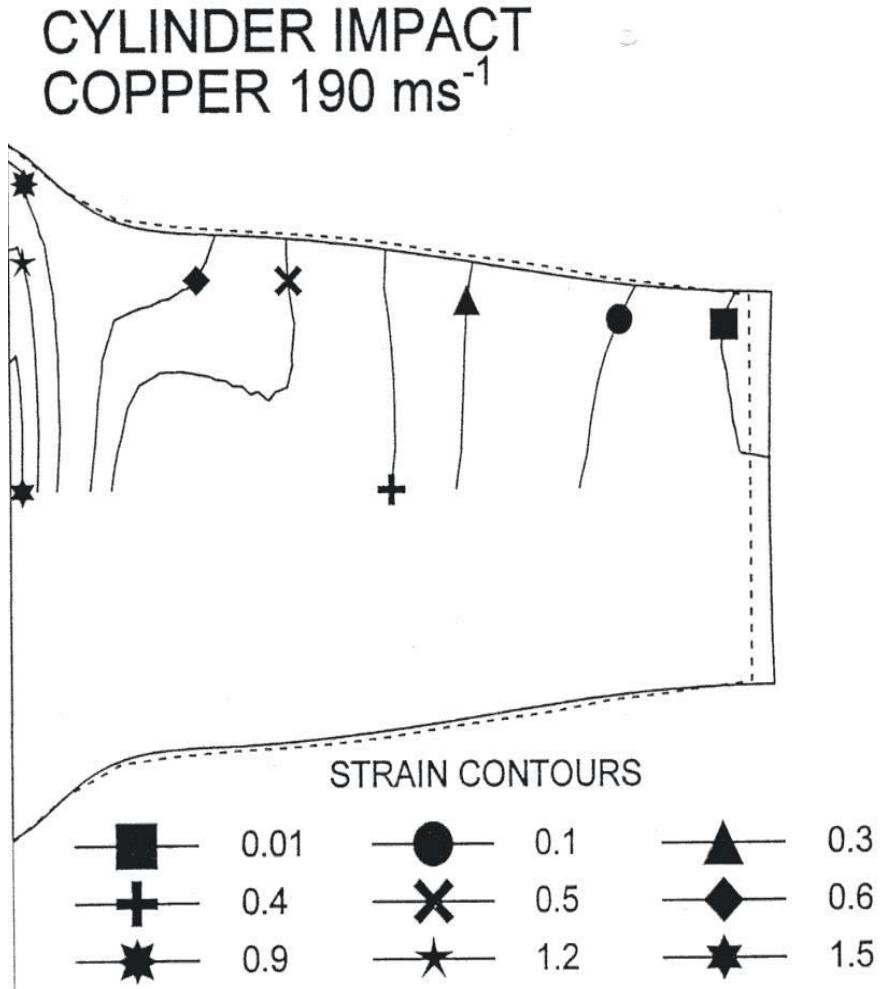


$$B_0 = 890 \text{ MPa}, \alpha_0 = 0.0028 \text{ K}^{-1}, \alpha_1 = 0.000115 \text{ K}^{-1}, (\varepsilon/\varepsilon_r) < 1.0, k_\varepsilon = 5 \text{ MPa.mm}^{1/2}, \sigma_G + k_\varepsilon t^{-1/2} = 65 \text{ MPa}$$

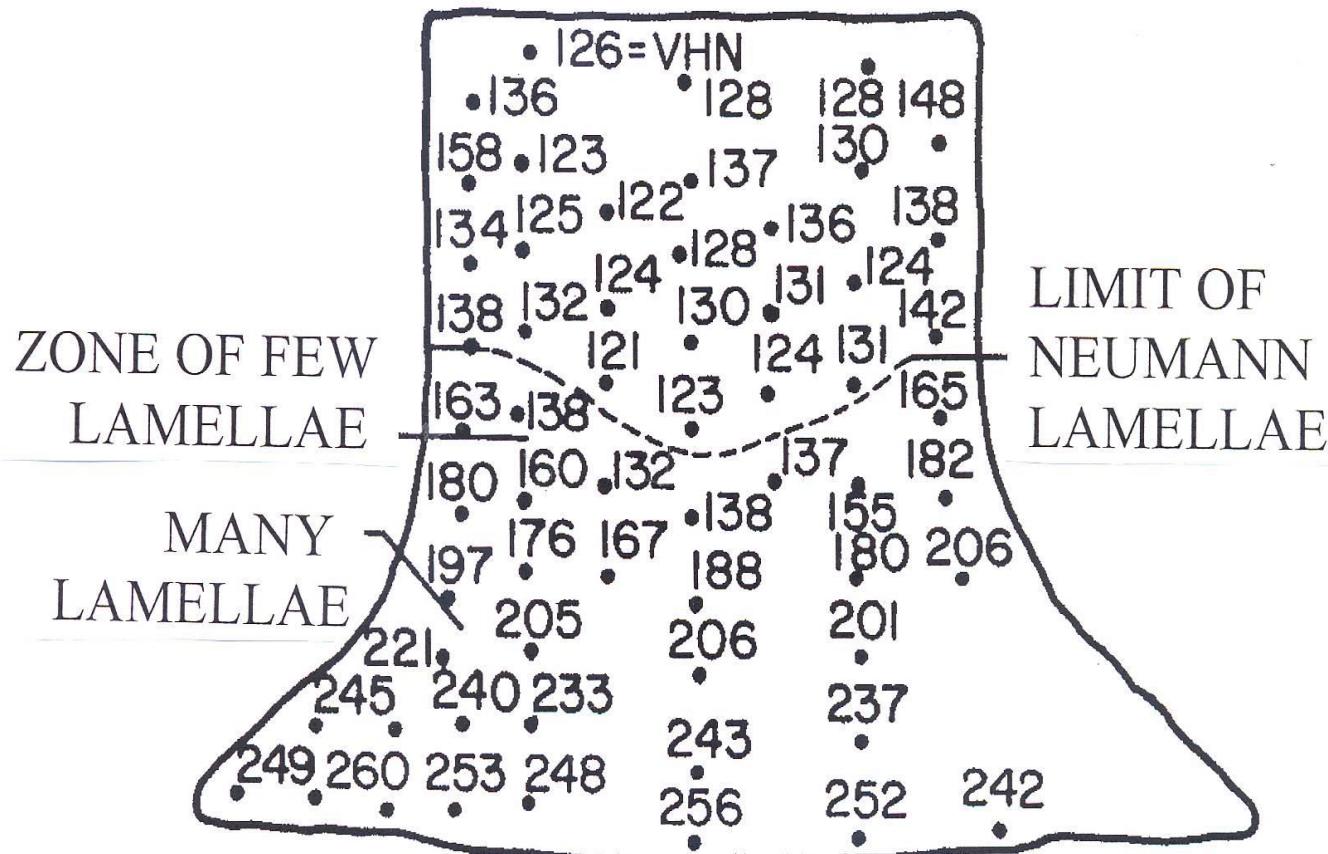
F.J. Zerilli and R.W. Armstrong, *J. Appl. Phys.* **61**, [5], 1816-1826 (1987);  
G.R. Johnson and W.H. Cook, *Eng. Fract. Mech.*, **21**, 31-48 (1985)

# J-C/Z-A fcc Taylor test result

$300 \leq T \leq \sim 600 \text{ K}$ ;  $0 \leq d\varepsilon/dt \leq \sim 10^5 \text{ s}^{-1}$

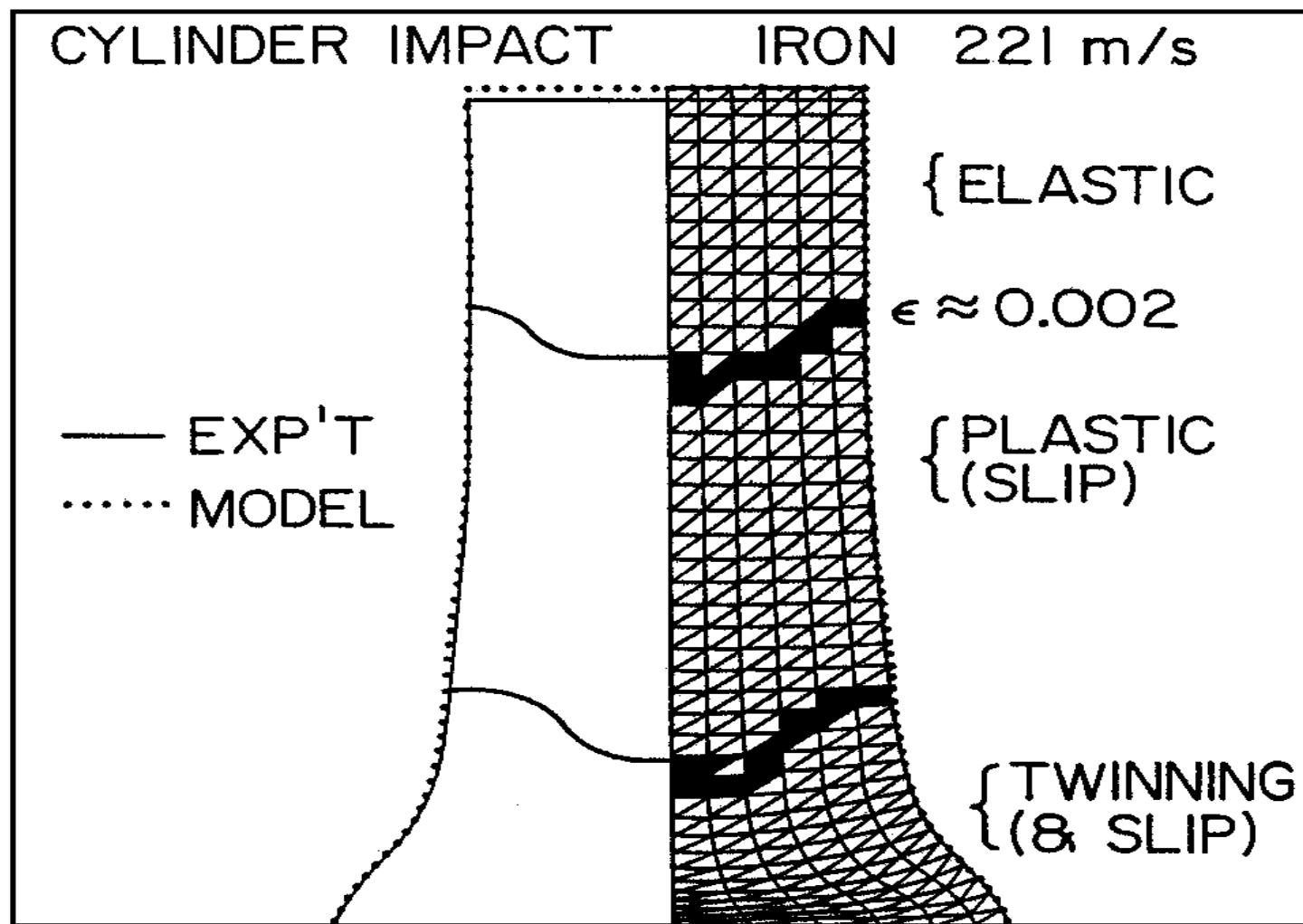


# Original Taylor cylinder test result on mild steel

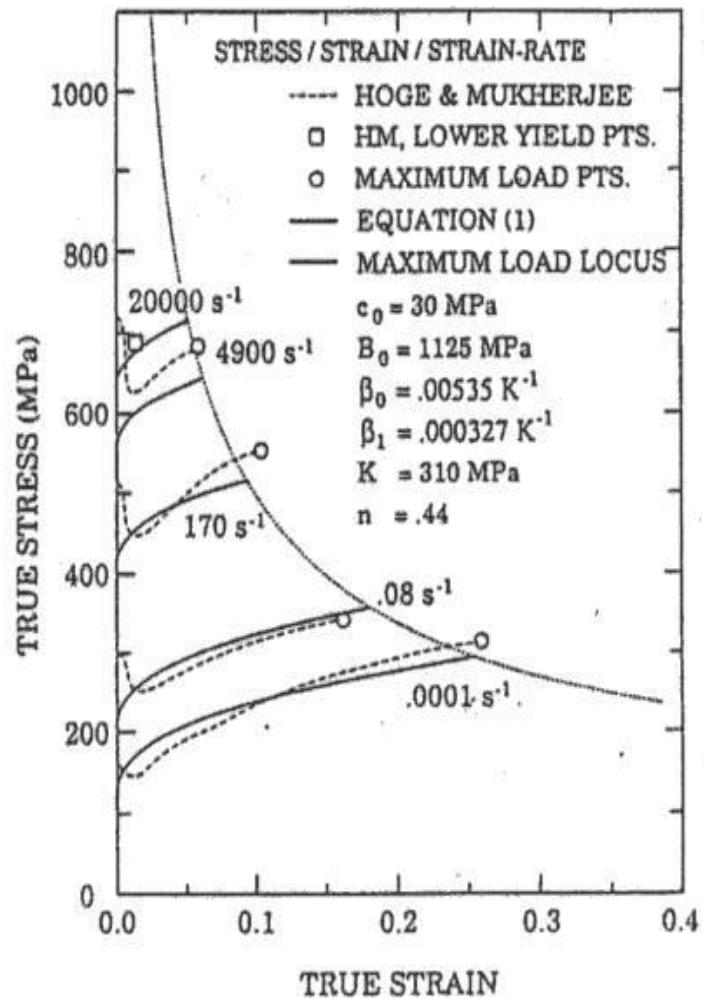
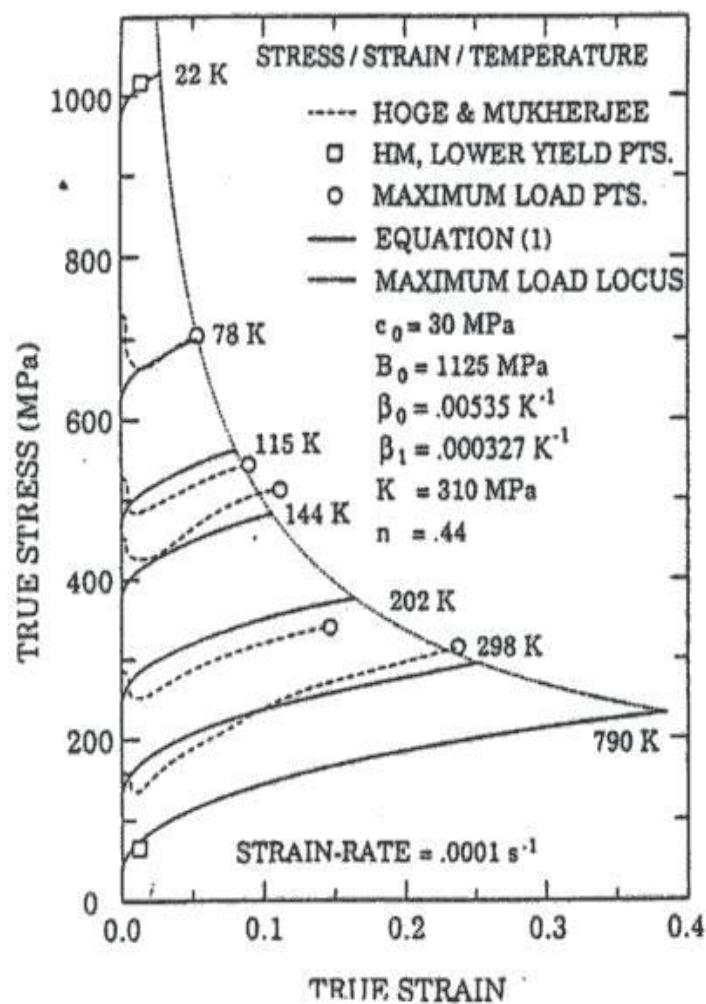


MILD STEEL CYLINDER. IMPACT VELOCITY  
338 m/s

# J-C/Z-A: twinning and slip in Armco iron

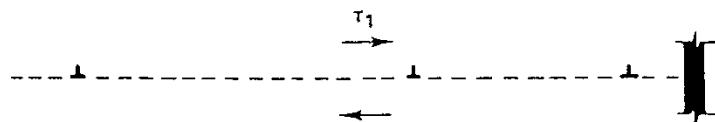


# Tantalum: $\sigma$ , $d\sigma/d\varepsilon$ over ( $T$ , $[d\varepsilon/dt]$ , $\varepsilon$ )



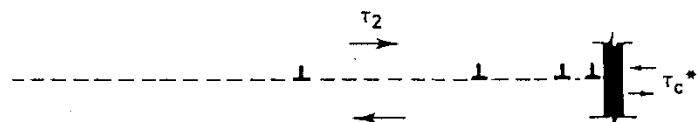
K.G. Hoge and A.K. Mukherjee, *J. Mater. Sci.*, **12**, 1666-1672 (1977); F.J. Zerilli, R.W. Armstrong, *J. Appl. Phys.*, **68**, 1580-1591 (1990); N.R. Barton, J.V. Bernier, R. Becker, A. Arsenlis, R. Cavallo, J. Marian, M. Rhee, H.-S. Park, B. Remington and R.T. Olson, *J. Appl. Phys.* **109**, 073501 (2011).

# Dislocation pile-up avalanche model for hot spots and shear banding



(a) isothermal stress build-up:  $n_1$  dislocations

$$\Delta T \leq \frac{k_s l^{1/2}}{16\pi} \left( \frac{2v}{c^* b K} \right)^{1/2} \quad \text{if } \left( \frac{2K}{c^* v b} \right) \leq 1.0$$



(b) critical stress concentration:  $n_2 \tau_2 = \tau_c^*$

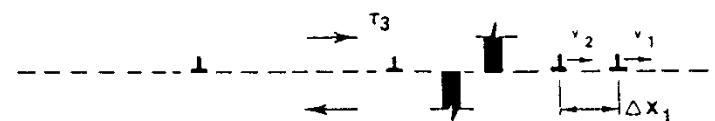
$$\dot{\gamma} = N b v,$$

$$v = v_0 \exp \left[ - \left( G_0 - \int_{\tau_{th}^0}^{\tau_{th}} bA^* d\tau_{th} \right) / kT \right],$$

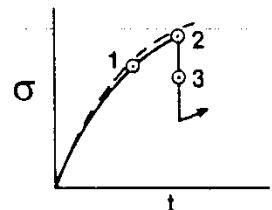
$$bA^* = W_0 / \tau_{th},$$

$$\tau_{th} \propto H_{50}^{1/n},$$

$$v = v_0 (H_{50}/H_{50}^0)^{W_0/nkT} \exp(-G_0/kT)$$



(c) adiabatic collapse-discontinuous load drop



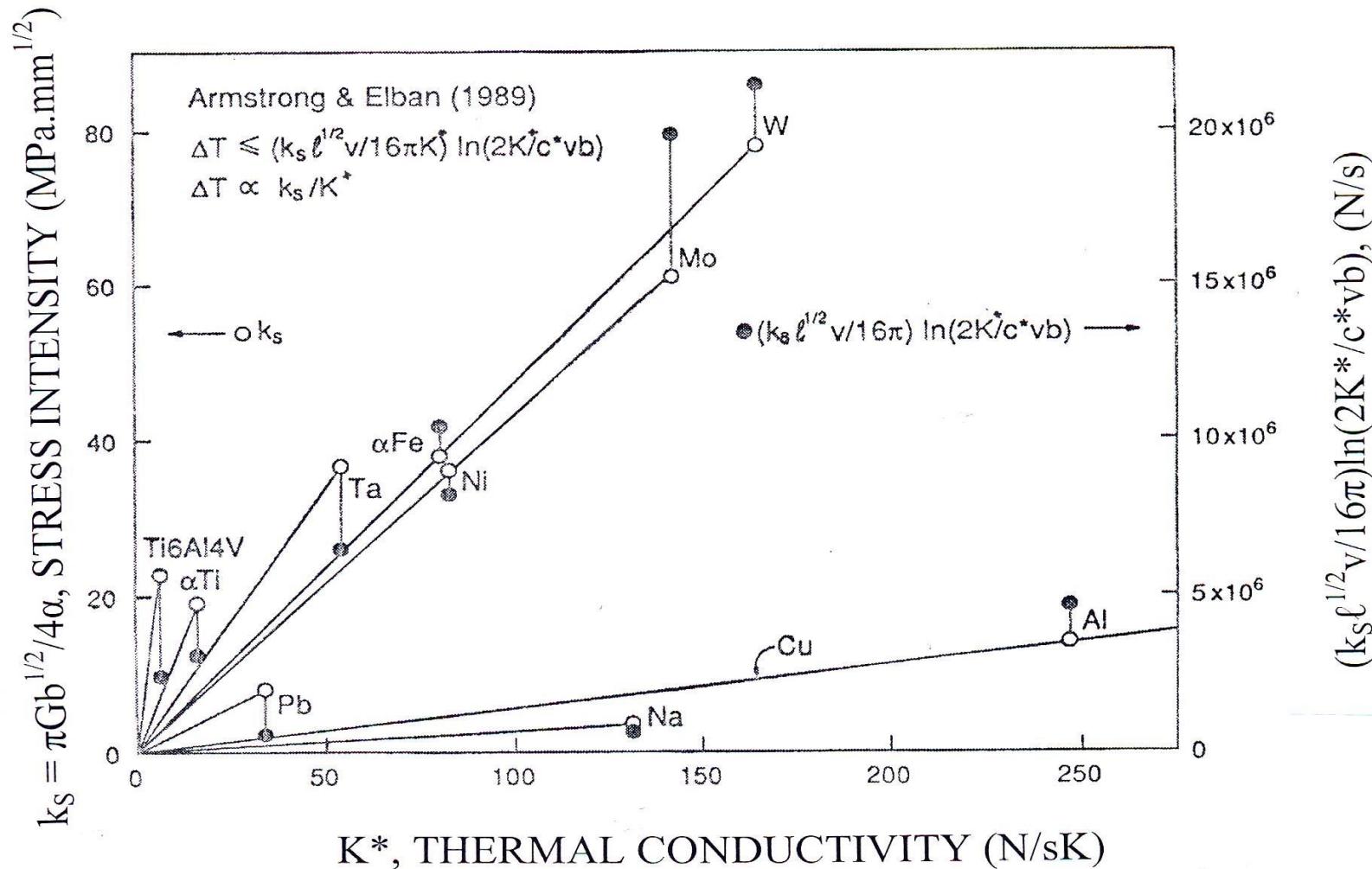
(d) pressure-time curve for  $\tau_1$ ,  $\tau_2$ , and  $\tau_3$

$$\begin{aligned} \log H_{50} &= \log H_{50}^* + (n/m^*) \log \{ [8\pi\sqrt{2}\Delta T \\ &\times K^{1/2} c^{*1/2} b^{1/2} / k_s v_0^{1/2} \\ &\times \exp(-G_0/2kT)] l^{-1/2} \}. \end{aligned}$$

$$m^* = W_0 / 2kT$$

R.W. Armstrong, C.S. Coffey and W.L. Elban, "Adiabatic heating at a dislocation pile-up avalanche", *Acta Metall.* **30**, 2111-2118 (1982)

# Shear band susceptibility: $k_s/K^*$



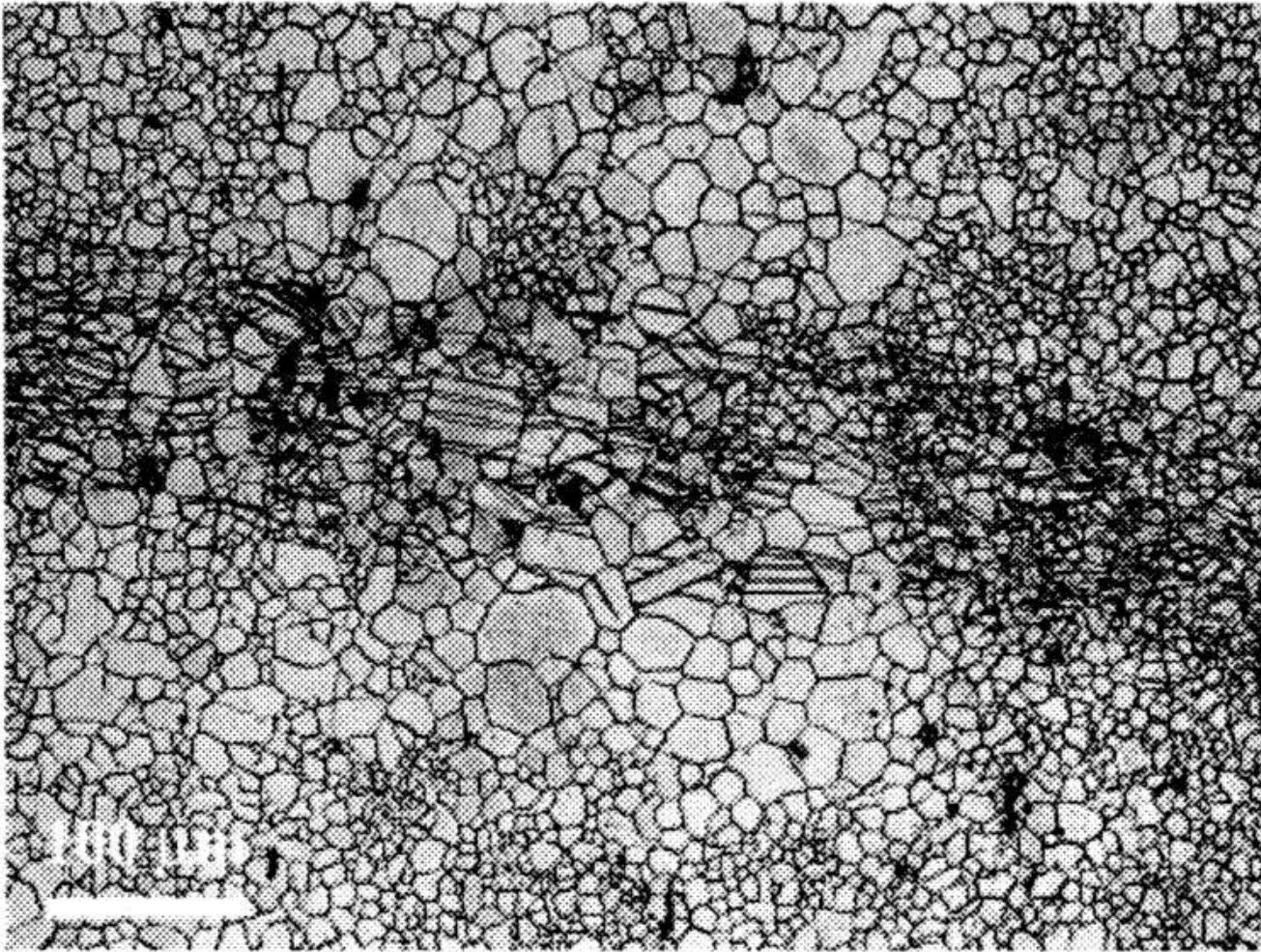
R.W. Armstrong and F.J. Zerilli, "Dislocation mechanics aspects of plastic instability and shear banding", In: ***Shear Instabilities and Viscoplasticity Theories***, edited by R.W. Armstrong, R.C. Batra, M.A. Meyers and T.W. Wright, *Mech. Mater.*, **17**, 319-327 (1994)

# Molten metal spray on Ti6Al4V shear plug



W. H. Holt, W. Mock, Jr., W.G. Soper, C.S. Coffey, V. Ramachandran and R.W. Armstrong, "Reverse-ballistic impact study of shear plug formation and displacement in Ti6Al4V alloy", *J. Appl. Phys.* **73**, 3753-3759 (1993).

# Lueders bands in Mg



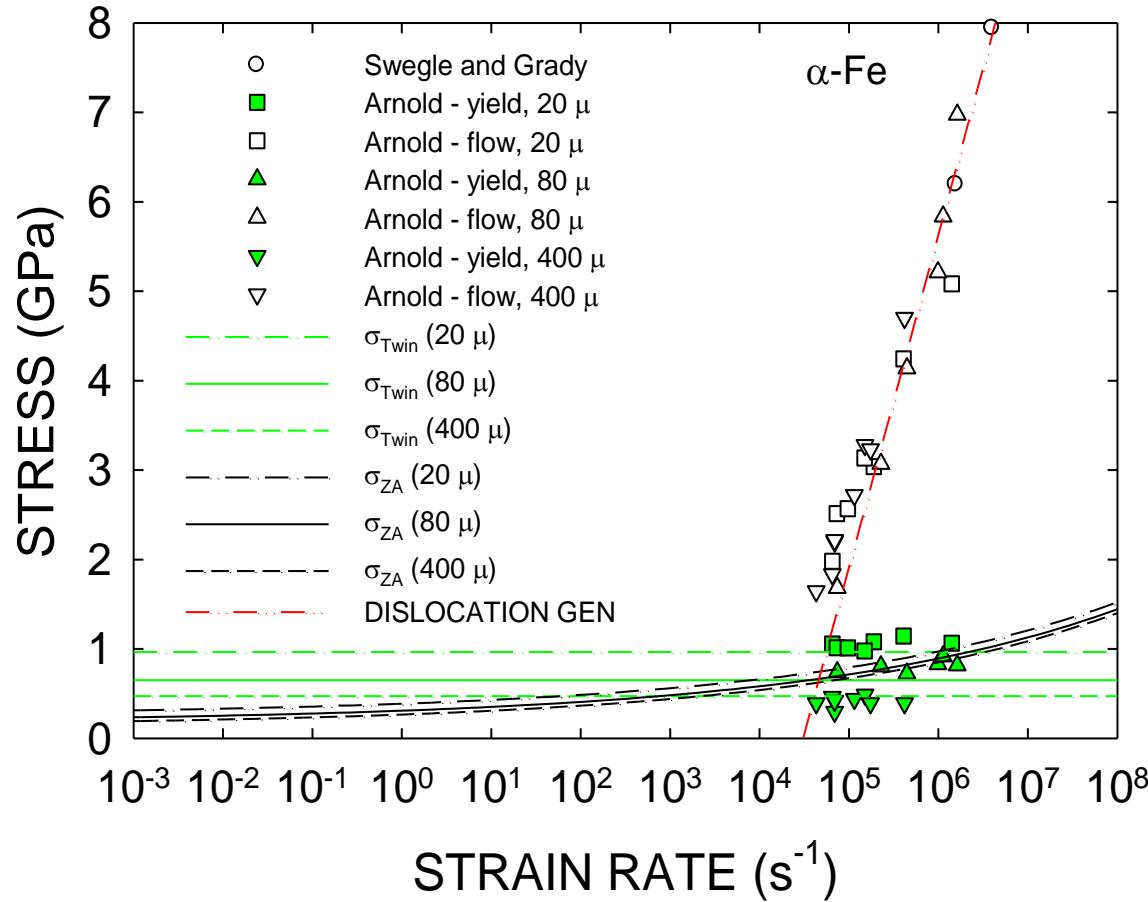
M.S. Tsai and C.P. Chang, "Grain size effect on deformation twinning in Mg-Al-Zn alloy (AZ31B)", Mater. Sci. Tech, **29**, 759-763 (2013); S.D. Antolovich and R.W. Armstrong, "Plastic strain localization in metals: Origins and consequences", Prog. Mater. Sci. **59**, 1-160 (2014)

# The progression

1. Beginning with the TASRA and leading to impact results obtained on copper and  $\alpha$ -iron materials, then to shear banding, a rather smooth progression occurs for higher strength levels being achieved at higher applied plastic strain rates.
2. A much stronger strain rate dependence is obtained for shock-induced plastic flow stress levels because of enhanced dislocation generation occurring along the propagating shock front.
3. A comparably high drag-controlled plastic flow stress occurs for isentropic compression experiments (ICEs) because the initially resident dislocation density must move at high speed.

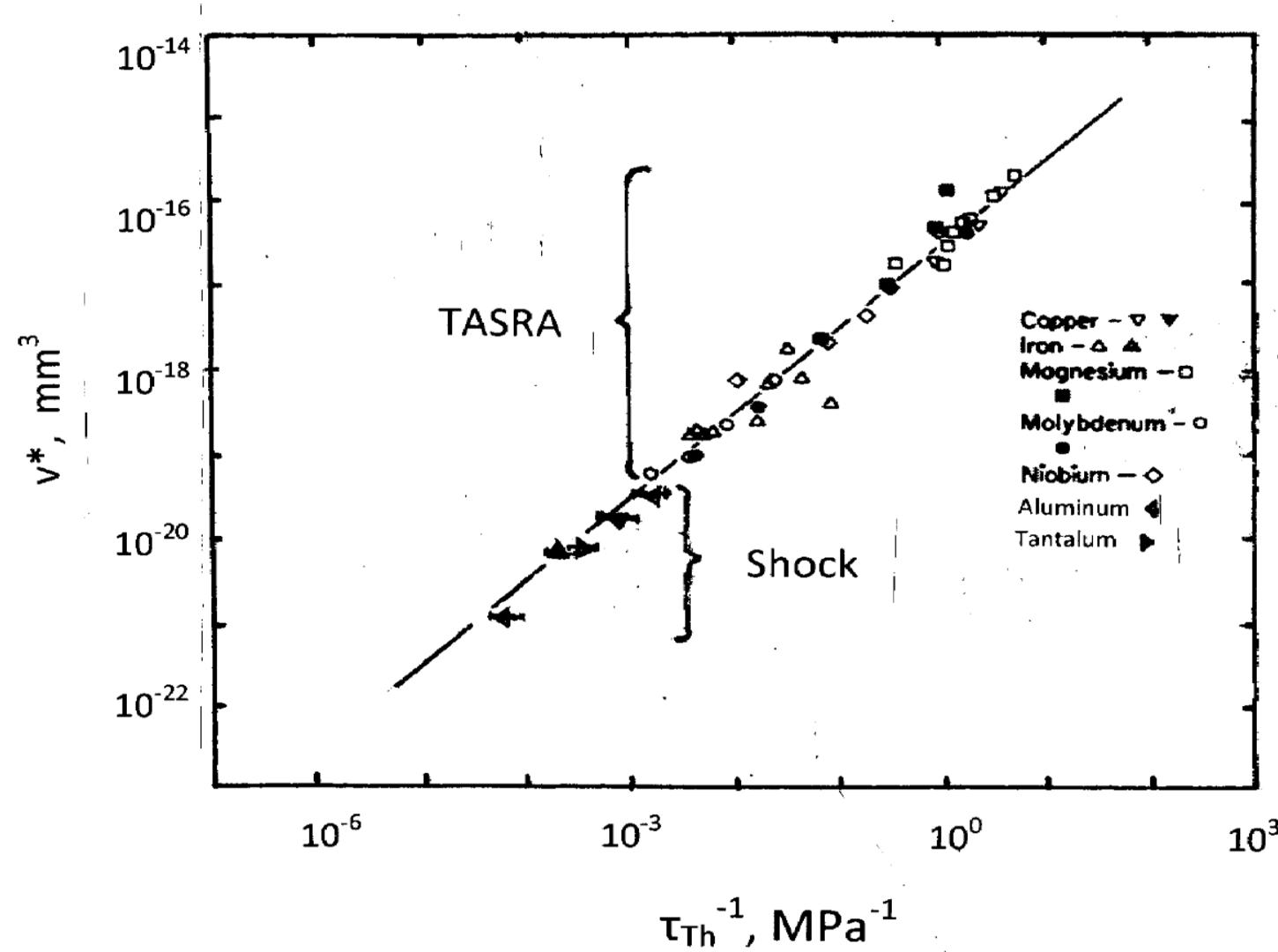
# Shock-induced plasticity in Armco iron

[twinning-determined HEL & grain-size-independent ( $d\varepsilon/dt$ )]



W. Arnold, *Dynamisches Werkstoffverhalten von Armco-Eisen Stosswellbelastung*, Fortschrittberichte VDI 5 (VDI Verlag GmbH, Duesseldorf, DE, 1992); R.W. Armstrong, W. Arnold and F.J. Zerilli, *J. Appl. Phys.*, **105**, 023511 (2009).

# TASRA correlation of shock measurements



# Dislocation drag control for (shockless) isentropic compression

The resident dislocation density is required to “carry the load”, and because  $\rho_N$  is low,  $v_N$  is so high as to be controlled by “drag”!

$$\sigma_{Th} = \{1 - [c(d\varepsilon/dt)/\beta_1 \sigma_{Th}]^{-\beta_1 T}\} [B \exp(-\beta T)]$$

in which

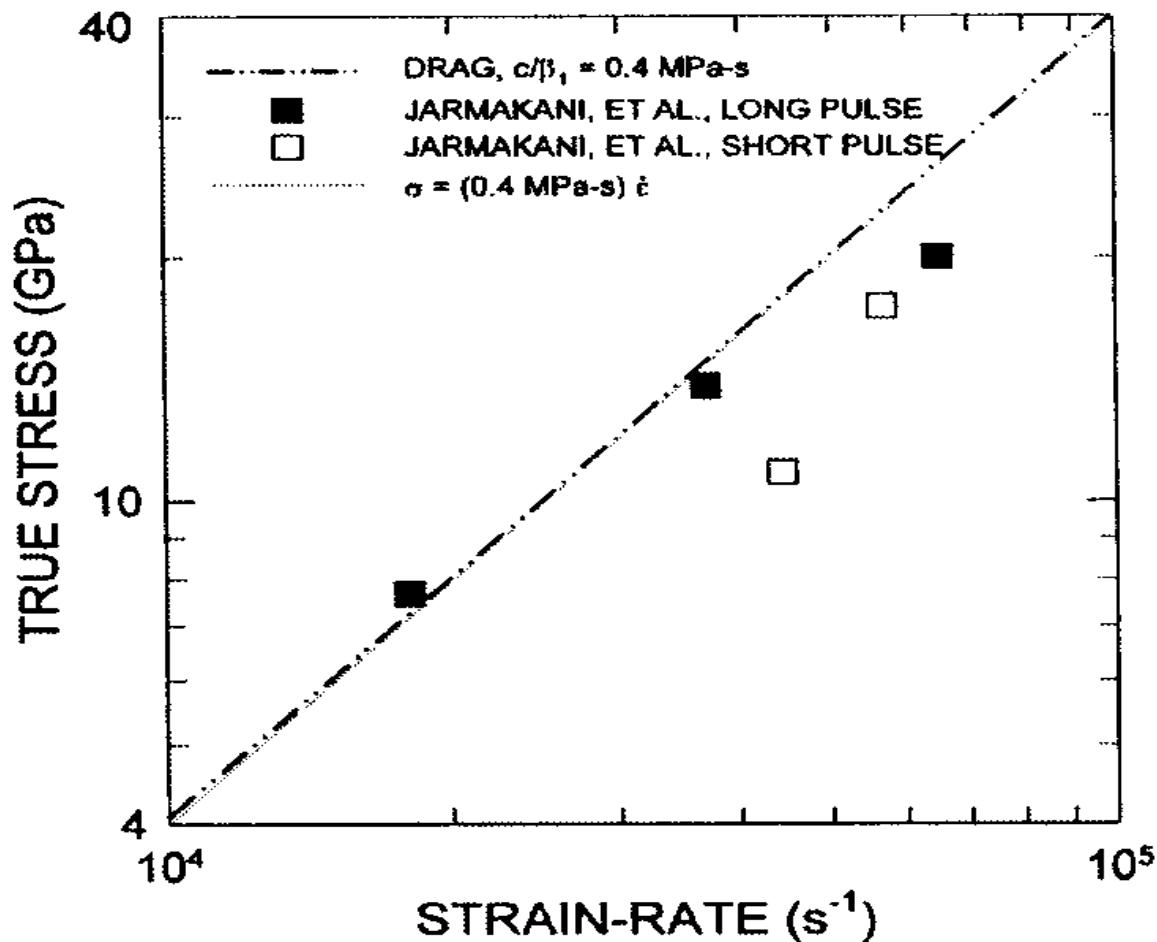
$$c = c_0 m^2 \beta_1 / \rho b^2 \quad \text{and} \quad b T_{Th} = c_0 v.$$

At limiting high ( $d\varepsilon/dt$ ):

$$\sigma_{Th} = (c_0 m^2 / \rho b^2) (d\varepsilon/dt)$$

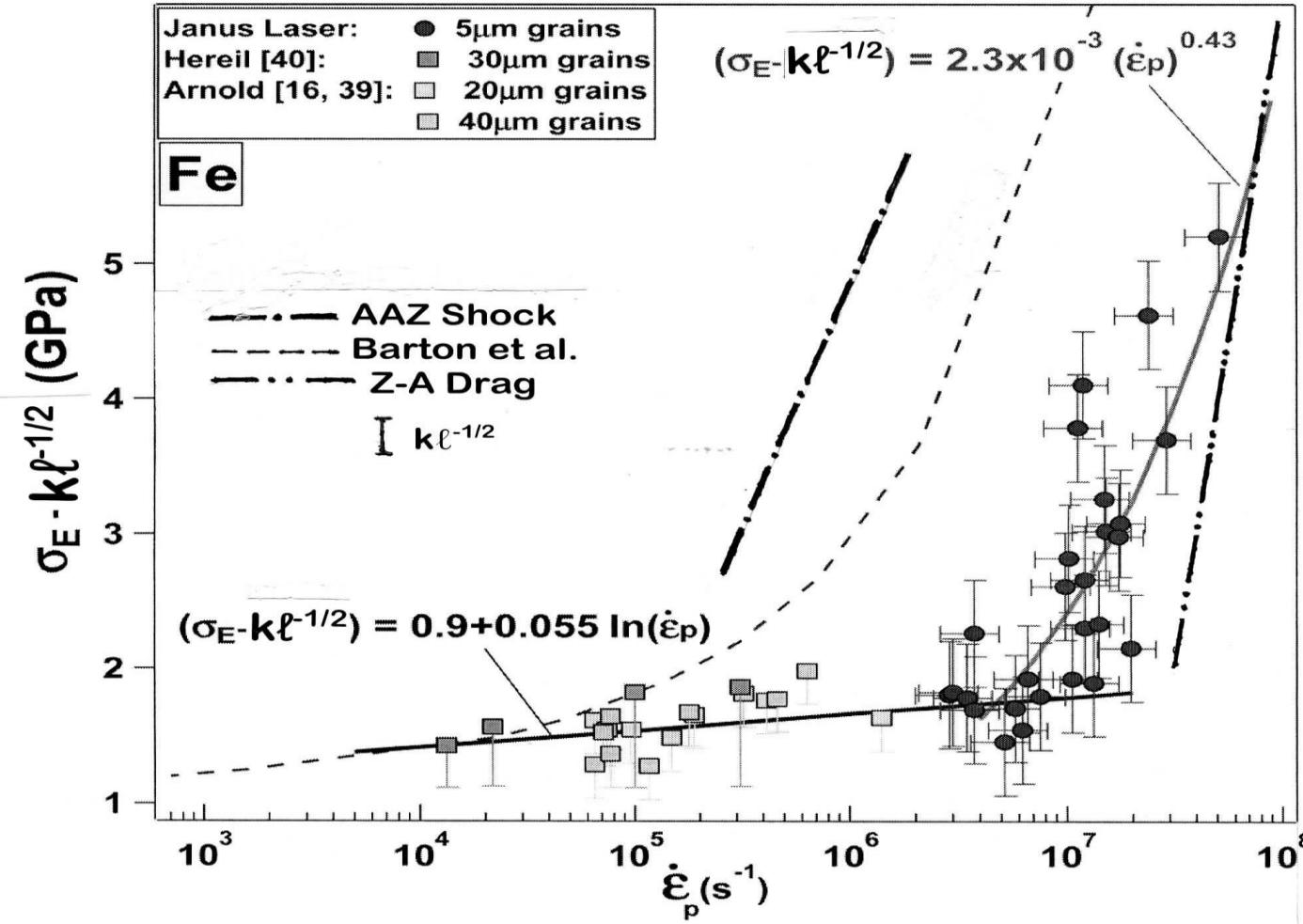
F.J. Zerilli and R.W. Armstrong, *Acta Mater.*, **40**, 1803-1808 (1992);  
R.W. Armstrong, W. Arnold and F.J. Zerilli, *J. Appl. Phys.* **105**, 023511 (2009)

# Drag control for (shockless) ICE on copper



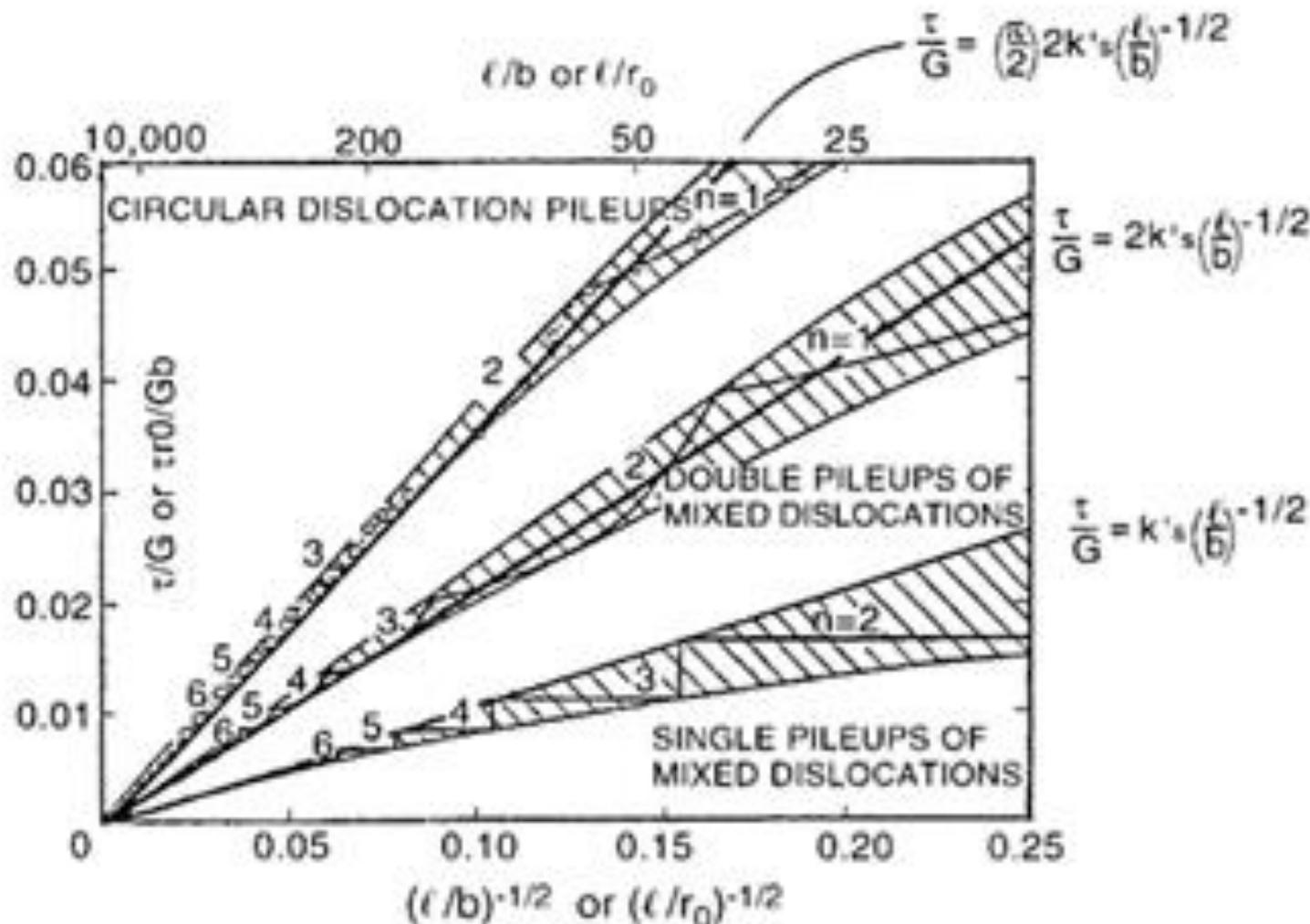
H. Jarmakani, J.M. McNaney, M.S. Schneider, D. Orlikowski et al., In: *Shock Compression of Condensed Matter - 2005*, M.D. Furnish, M. Elert, T.P. Russell and C.T. White, eds., (Amer. Inst. Phys., Melville, NY, 2006) **CP845**, Part 2, pp. 1319-1322; R.W. Armstrong, W. Arnold and F.J. Zerilli, *J. Appl. Phys.*, **105** 023511 (2009)

# SHPB, shock and ICE results on ARMCO iron



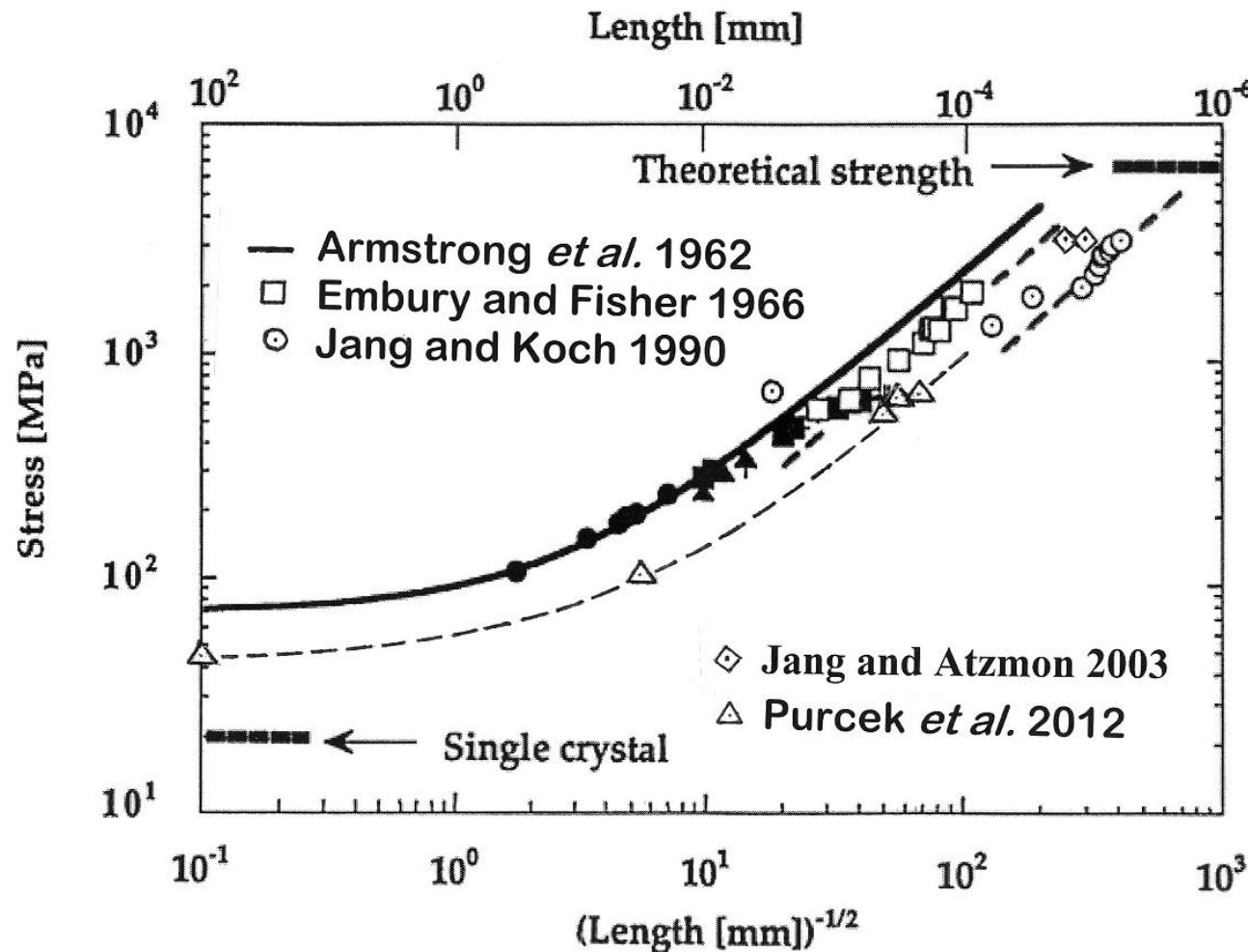
R.F. Smith, J.H. Eggert, R.E. Rudd, D.C. Swift, C.A. Bolme and G.W. Collins, *J. Appl. Phys.*, **110**, 123515 (2011), see Fig. 9; N.R. Barton, J.V. Bernier, R. Becker, A. Arsenlis, R. Cavallo, J. Marian, M. Rhee, H.-S. Park, B. Remington and R.T. Olson, *J. Appl. Phys.* **109**, 073501 (2011); R.W. Armstrong, W. Arnold and F.J. Zerilli, *J. Appl. Phys.* **105**, 023511 (2009).

# Hall-Petch/Griffith connection for $k$ being described as a stress intensity



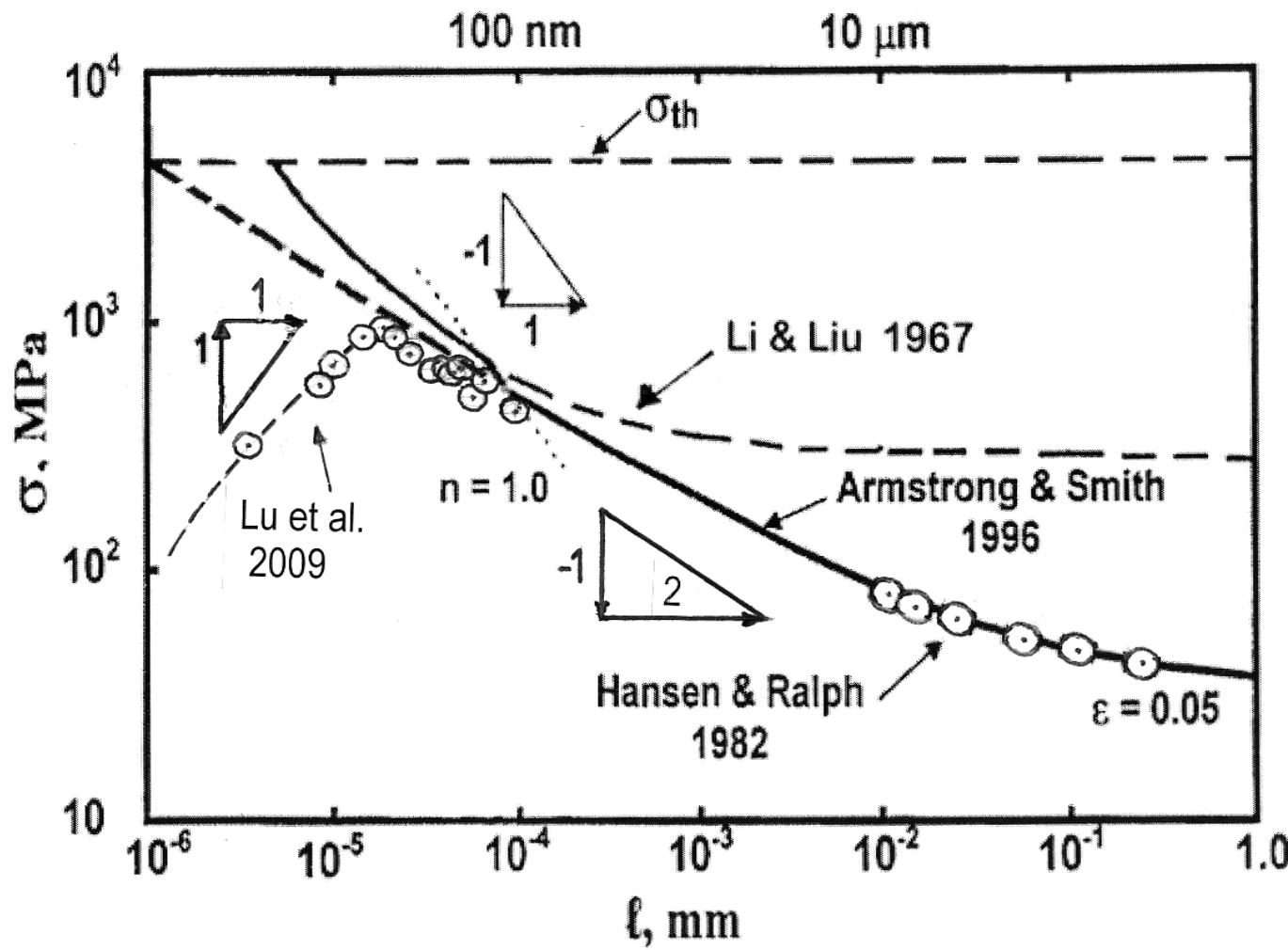
R.W. Armstrong, "Dislocation pile-ups: from {110} cracking in MgO to model strength evaluations", *Mater. Sci. Eng. A*, **409**, 24-31 (2005)

# The H-P dependence for iron and steel on a log/log basis



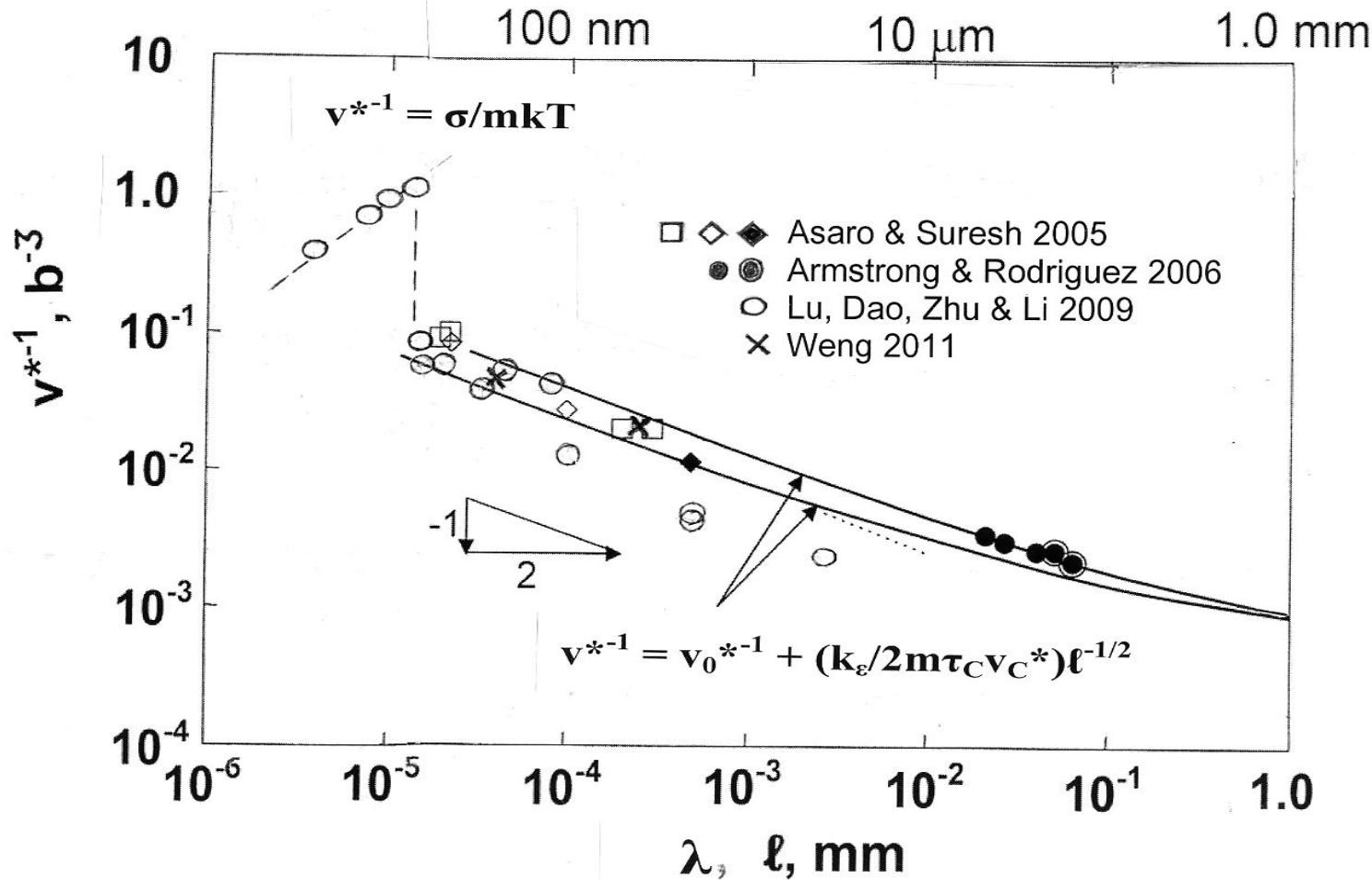
R.W. Armstrong, “Hall-Petch analysis: Past to present nano-scale connections”, in **Strength of Fine Grained Materials – 60 Years of Hall-Petch**, (Japan) Mater. Trans. (2014)

# Limiting H-P relation for one dislocation loop and H-P application for twin boundaries in copper



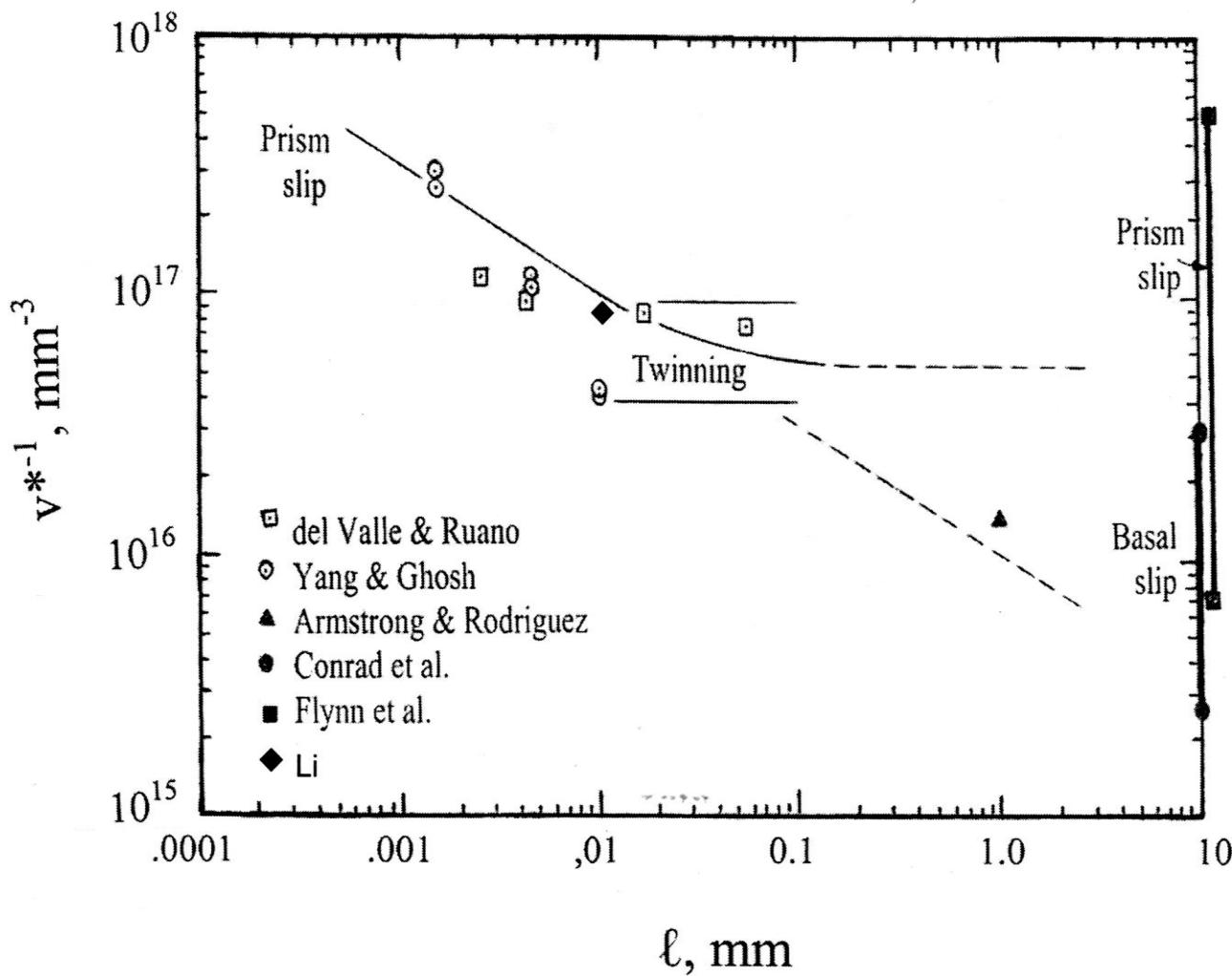
R.W. Armstrong, "60 Years of Hall-Petch: Past to Present Nano-Scale Connections", (*Japan*) Mater. Trans., 55, [1], 2-12 (2014).

# H-P description of Cu and Ni strain rate sensitivity,

$$v^* = kT \{ \partial \ln [dy/dt] / \partial T \}_T$$


R.W. Armstrong and P. Rodriguez, *Philos. Mag.* **86**, 5787-5796 (2006); R.W. Armstrong, In: **Strength of Fine Grained Materials – 60 Years of Hall-Petch**, *Mater. Trans.* **55**, 2-12 (2014).

# Log/log grain size dependence of $v^*$ for Mg



R.W. Armstrong, In: *Nanometals – Status and Perspective*, 33<sup>rd</sup> Risoe Intern. Symp. On Mater. Sci. (Technical University of Denmark, Roskilde, DK, 2012) pp. 181-199;  
Q. Li, *Mater. Sci. Eng. A*, **540**, 130-134 (2012)

# SUMMARY

1. Dislocation dynamics results were described on the basis of the thermal activation – strain rate analysis and the Z-A equations (1973, 1987, 1999).
2. Dislocation pile-up avalanches are proposed to provide a fundamental explanation of shear banding behavior (1981/2, 1993/4).
3. Modeling of shock-induced deformations places emphasis on dislocation generations (1992, 2009).
4. Shockless Isentropic Compression Experiments are explained in terms of dislocation drag (ICEs, 2006/7).
5. Hall-Petch results were demonstrated for nanopoly-crystalline materials (1962, 1969/70, 2006).